



Crack growth rate prediction based on damage accumulation functions for creep-fatigue interaction

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ABSTRACT. The present study is concerned with formulation of a model for the creep–fatigue crack growth rate prediction on the base of fracture damage zone concepts. It is supposed that crack growth rate can be determined by integration of damage accumulation rate equations into the fracture process zone for low-cycle fatigue and creep loading independently. In the case of low-cycle fatigue loading the damage accumulation function proposed by Ye and Wang was used as well as a classical Kachanov-Rabotnov power law was employment for the creep damage accumulation characterization. Fracture process zone size is calculated on the base of the nonlinear stress intensity factors concept proposed by Shlyannikov. The background for the proposed general model of crack growth rate under creep and fatigue interaction is given in order to comparison with the experimental data. Experimental study of crack growth rate under creep and fatigue interaction is performed for compact tension specimen made from 20CrMoV5. Crack growth rate carried out at the elevated temperature of 550°C according to ASTM E2760 standard. The predictions of the crack growth rate were compared with the experimental data for the 20CrMoV5 steel obtained at an elevated temperature, and the agreement was found to be satisfactory.

KEYWORDS. Creep-fatigue interaction; Crack growth rate prediction; Nonlinear stress intensity factors.



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INTRODUCTION

Structures in the aviation, space flight, power generation, and petrochemical industries are typically exposed to elevated temperatures. Damage accumulation and growth considerations at creep and fatigue owing to changes in the material microstructure, void nucleation, interaction, and growth on the grain boundaries are important in the design and operation of such components in order to ensure structural integrity. It is well known that the creep damage accumulation from the damage caused by cyclic loading is different. At the moment, there are many models in literature allow to take into account the damage accumulation, both during creep and cyclic loading.



The Hayhurst model [1] one of the generalized models for the creep damage accumulation. The main disadvantage of this model is necessity to use six parameters as material properties in order to describe the damage accumulation behavior. Moreover, the parameters determination methods are mainly indirect. According to this, the Rabotnov-Kachanov model is used due to the direct method of material constants determination.

The Lemaître model [2] is the one of the most frequently used for the fatigue damage accumulation behavior. In similar phenomenological models using the tensor damage parameter leads to considerable difficulties for material constants determination. For the different creep and fatigue laws if the stress state is unchanged the scalar, vector, and tensor damage parameters give almost identical results. In this regard, models with a scalar damage parameter as a simplified alternative can be considered [3-6].

The main problem of the residual life prediction under the creep-fatigue loading conditions is the complex character of their interaction [7-10]. Statistical approximations of experimental results by the polynomial functions are most common approach of the creep-fatigue interaction behavior. In present study, continuum damage mechanics is applied to assess the creep-fatigue interaction and the crack growth rate prediction, considering the complete sequence of formulations.

FATIGUE DAMAGE ACCUMULATION RATE

The variable critical distance is one of important parameter in modern fracture mechanics, denoted to as the fracture process zone size. A general assumption regarding the distance criterion under elastic-plastic and creep loading conditions is that a crack increment occurs when the fracture resistance parameter (stress, strain, or energy) reaches a critical value at a characteristic distance from the crack tip.

In creep-fatigue interaction both the fatigue and the creep crack growth rate can be obtained from damage accumulation rule independently. The crack size is assumed to increase when the local strain energy density at the crack tip reaches the critical value. Hence, a material point initially at a characteristic distance, r_c , ahead of the crack tip, $r = 0$, moves to the tip both in time Δt after an increment in the crack growth for creep and in number of cycles ΔN , in case of fatigue. It is assumed that the crack length increment equal to the fracture process zone size, $\Delta a = r_f$, and the crack growth rate, \dot{a} , becomes [11-12]:

$$\dot{a}_f = \int_0^{r_f} \dot{\omega}(r_f) dr_f \quad (1)$$

where $\dot{\omega}(r_f)$ is the damage accumulation rate.

Eqn. (1) for creep conditions was proposed initially, however the analogy between the creep time Δt and the number of cycles ΔN makes possible to use this dependence for low cycle fatigue. Usually, the one of the main characteristic in models with a scalar damage parameter is the number of cycles before fracture N_f at the stress amplitude σ_a . In the simplest case of uniaxial tension-compression, the fatigue life can be obtained from:

$$2N_f = \left(\frac{\sigma_a}{\sigma'_{f_f}} \right)^c \quad (2)$$

where σ'_{f_f} and c are fatigue constants.

The damage parameter at constant stress amplitude according to [13]

$$\omega_f = -\frac{\omega_{f_{N-1}}}{\ln N_f} \ln \left[1 - \frac{N}{N_f} \right] \quad (3)$$

where $\omega_{f_{N-1}}$ - the critical value of the damage parameter in the penultimate cycle and it is associated with static toughness U_{T_0} by the following relationship:



$$\omega_{f_{N-1}} = 1 - \frac{\sigma_a^2}{2EU_{T_0}} \quad (4)$$

where E - Young's modulus.

The fatigue damage accumulation rate can be obtained from the equation (3)

$$\dot{\omega}_f = \frac{d\omega_f}{dN} = \frac{\omega_{f_{N-1}}}{\left[(N_f - N) \ln N_f \right]} \quad (5)$$

FATIGUE CRACK GROWTH RATE

For Ramberg-Osgood hardening law [26] the critical value of a strain energy density W_c can be obtained from true stress-strain diagram:

$$W_c = \int_0^{\sigma_f} \left[\frac{\sigma}{E} + \alpha \frac{\sigma_y}{E} \left(\frac{\sigma}{\sigma_y} \right)^{n_p} \right] d\sigma \quad (6)$$

where α - is the strain hardening coefficient, n_p - is the strain hardening exponent, σ_y - is the yield stress, σ_f - ultimate tensile stress.

According to the classical Hutchinson-Rosengren-Rice model in the zone where fully plastic singularity is dominated the strain energy density W can be found from following expression [14]:

$$W = \alpha \frac{n_p}{n_p + 1} K_p^{n_p+1} \tilde{\sigma}_e^{n_p+1} \frac{\sigma_y^2}{E} \frac{(1-\nu^2)}{\bar{r}} \quad (7)$$

where ν - is Poisson's ratio, $\bar{r} = r/a$ - crack tip distance, a - crack length, $\tilde{\sigma}_e$ - dimensionless equivalent Mises stresses which depend only on a polar coordinate θ and normalized by $\tilde{\sigma}_e(\theta)|_{\theta=0}^{\max} = 1$. The plastic stress intensity factor K_p in small-scale and extensive yielding conditions in Eqn.(7) can be expressed directly using Rice's J -integral [15]. That is

$$\bar{K}_p = \left[\frac{J E}{\alpha_p \sigma_y^2 I_p L} \right]^{\frac{1}{n_p+1}} \quad (8)$$

where

$$\begin{aligned} J = & \frac{r \sigma_y^2}{E} \int_{-\pi}^{+\pi} \left[\frac{1+\nu}{3} \bar{\sigma}_e^2 + \frac{1-2\nu}{6} \bar{\sigma}_{kk}^2 + \frac{\alpha_p n}{n+1} \bar{\sigma}_e^{n+1} \right] \cos \theta d\theta - \\ & - r \sigma_y \int_{-\pi}^{+\pi} \left(\bar{\sigma}_r \frac{\partial \bar{u}_r}{\partial r} + \bar{\sigma}_{r\theta} \frac{\partial \bar{u}_\theta}{\partial r} \right) \cos \theta d\theta + \sigma_y \int_{-\pi}^{+\pi} \left(\bar{\sigma}_{rr} \frac{\partial \bar{u}_r}{\partial \theta} + \bar{\sigma}_{r\theta} \frac{\partial \bar{u}_\theta}{\partial \theta} \right) \sin \theta d\theta \end{aligned} \quad (9)$$

In Eqns. (8-9), \bar{u}_i are crack tip dimensionless displacement components, L is cracked body characteristic size, in our case is the crack length $L=a$, and I_p is the numerical constant of the crack-tip stress-strain field, which should be



determined for the cracked body with a finite size. In the case of full three-dimensional cracked body, the I_p -factor in Eqn. (8) changes not only with the strain hardening exponent n_p but also with the position along the crack front. In Eqn. (9), the stress tensor and invariant are both normalized by the yield stress: $\bar{\sigma}_{ij} = \sigma_{ij}/\sigma_y$ and $\bar{\sigma}_{kk} = \sigma_{kk}/\sigma_y$. More detail in determining the I_p factor for different elastic-plastic cracked body configurations are given by Refs. [15-18].

In a first approximation the crack growth as a static cracks sequence is considered. Thus, the path independent J-integral can be obtained from:

$$J = \frac{K_1^2(1-\nu^2)}{E} \quad (10)$$

Substitution Eqns. (6) and (8) into (7) give a possibility to obtain the critical distance where the strain energy density reaches a critical value:

$$\bar{r}_f = \alpha \frac{n}{n+1} K_p^{n_p+1} \tilde{\sigma}_e^{n_p+1} \frac{\sigma_y^2}{E} \frac{(1-\nu^2)}{W_c} \quad (11)$$

The fatigue crack growth rate can be calculated by substitution Eqns. (5) and (11) into (1) directly, it leads to the following relationship:

$$\left(\frac{da}{dN} \right)_{fatigue} = \int_0^{r_f} \dot{\phi}_f dr = \alpha \sigma_y^2 \tilde{\sigma}_e^{n_p+1} K_p^{n_p+1} \frac{\alpha n_p}{n_p+1} \left[\frac{\sigma_y^2 \tilde{\sigma}_e^2 r_f^{n_p+1} K_p^2}{2 E U_{t0}} - 1 \right] \frac{(1-\nu^2)}{E W_c \ln(N_f)(N-N_f)} \quad (12)$$

CREEP DAMAGE ACCUMULATION

In this study the classical Kachanov-Rabotnov power law is used for the creep damage accumulation description. According to this model, the strain rate during creep is [1-2]:

$$\dot{\varepsilon}_e = B \left(\frac{\sigma_e}{1-\omega_\sigma} \right)^{n_\sigma} \quad (13)$$

and the creep damage accumulation rate is:

$$\omega_\sigma = \frac{d\omega_\sigma}{dt} = D \left(\frac{\sigma_e}{1-\omega_\sigma} \right)^m \quad (14)$$

where B and n_σ are material constants of the Norton power law constitutive equation, D and m – are material properties. The damage variable ω indicate the measure of creep damage with $\omega = 0$ denoting the undamaged state and $\omega = 1$ the fully damaged state. The creep damage increment $\Delta\omega_\sigma$ at time Δt can be obtained by integrating the expression (14):

$$\Delta\omega_\sigma = 1 - \left[\left(-D \sigma'' \Delta t - C_3 \right) (m+1) \right]^{\frac{1}{m+1}}. \quad (15)$$

where C_3 - is integration constant. This constant can be determined from the initial state.



$$C_3 = -\frac{(1-\omega_0)^{m+1}}{m+1} - D\sigma_e^m t_0 \quad (16)$$

Based on the fact that the damage parameter is equal to one at failure $\omega = 1$ the fracture time at the known stress level can be calculated directly:

$$t_{cr} = \frac{1}{(m+1)D\sigma_e^m} \quad (17)$$

CREEP CRACK GROWTH RATE

For elastic-nonlinear-viscous material behavior, the stress, strain and displacement rate fields can be used in order to account for a creep stress intensity factor K_{cr} , which is amplitude of singularity. For extensive creep conditions the relation between the C -integral and creep stress intensity factor is introduced by the authors [19] in the form:

$$\bar{K}_{cr} = \frac{1}{\sigma_{ref}} \left(\frac{C^*}{BI_{cr}L} \right)^{\frac{1}{n_{cr}+1}} \quad (18)$$

$$C^* = \int_{-\pi}^{+\pi} Br \frac{n_{cr}}{n_{cr}+1} \sigma_e^{n_{cr}+1} \cos \theta d\theta - \int_{-\pi}^{+\pi} \left[\sin \theta \left(\sigma_r \frac{\partial \dot{u}_r}{\partial \theta} + \sigma_{r\theta} \frac{\partial \dot{u}_\theta}{\partial \theta} \right) - r \cos \theta \left(\sigma_r \frac{\partial \dot{u}_r}{\partial r} + \sigma_{r\theta} \frac{\partial \dot{u}_\theta}{\partial r} \right) \right] d\theta \quad (19)$$

where σ_{ref} - is reference stress, K_{cr} is amplitude of singularity in the form of creep stress intensity factor, C^* is the C -integral, \dot{u}_i are displacement rate angular functions, σ_{ij} are stress components. It should be noted that the I_{cr} - integral values are determined similar to I_p and can be determined directly from the finite element analysis by distribution of the displacement rate functions, \dot{u}_i , and dimensionless angular stress functions, $\tilde{\sigma}_{ij}$, [9]. More detail in determining the I_{cr} - integral for different creeping cracked body geometries are given by Refs. [12, 19-22]. For a static crack in the outer region of the small-scale creep zone, the elastic crack-tip field still dominates. In this case, the expression of C -integral has a simple form [5]:

$$C(t) = \frac{K_1^2(1-\nu^2)}{E(n_{cr}+1)t_{cr}} \quad (20)$$

The elastic stress intensity factor for a compact tension specimen is [23]:

$$K_1 = \frac{F}{b\sqrt{w}} Y_1 \quad (21)$$

where F - is applied load, Y_1 - is geometry correction function:

$$Y_1 = \frac{(2+a/w)}{(1-a/w)^{3/2}} (0.886 + 4.64a/w - 13.31 \left(\frac{a}{w} \right)^2 + 14.72 \left(\frac{a}{w} \right)^3 - 5.6 \left(\frac{a}{w} \right)^4) \quad (22)$$



The creep crack growth rate was determined based on the assumption that the crack propagation energy is equivalent to the damage initiation energy in infinite plate. Thus, the infinite plate remote stresses were calculated from the following expression:

$$\sigma_{ref} = \frac{K_1}{\sqrt{\pi a}} \quad (23)$$

Substitution Eqn.(23) into (17) give a possibility to calculate a reference failure time. A strain energy rate density parameter, \dot{W} , is employed by the author [20] to obtain the critical distance, r_{cr} , in the crack tip vicinity under creep conditions. If σ_{fc} referred the maximum value of the creep-rupture strength at a specified temperature and creep time, then the creep crack tip critical distance $r = r_{cr}$ is given by expression:

$$r_{cr} = a \left(\frac{\bar{K}_{cr}}{\bar{\sigma}_{fc}} \right) \tilde{\sigma}_e^{n_{cr}+1}. \quad (24)$$

where $\tilde{\sigma}_e$ is the dimensionless equivalent von Mises stress. Substitution of Eqns. (14) and (24) into Eqn. (1) with take into account that $r_{cr} = a\bar{r}_{cr}$, and integrating leads to the equation for creep crack growth in terms of C^* -integral and creep stress intensity factor \bar{K}_{cr} :

$$\left(\frac{da}{dt} \right)_{creep} = \int_0^{r_{cr}} \omega_{cr} dr = aD \left(\frac{\tilde{\sigma}_e}{\bar{\sigma}_{fc}} K_{cr} \right)^{n+1} \left(\frac{\sigma_{fc}}{K_{cr}^2} \right)^m \left[(1 - \omega_0)^{m+1} + D(m+1) \left(\frac{\sigma_{fc}}{K_{cr}^2} \right)^m (t_0 - t) \right]^{\frac{m}{m+1}} \quad (25)$$

CREEP-FATIGUE INTERACTION

The general simple superposition lifetime prediction models for creep–fatigue interaction can be differentiated as those that account for the hold-time effects and those that are employed for continuous cyclic scenarios. Therefore, we will use a linear summation law for creep-fatigue crack growth rate prediction [22, 24]

$$\frac{da}{dN} = \left(\frac{da}{dN} \right)_{fatigue} + \left(\frac{da}{dN} \right)_{creep} = \left(\frac{da}{dN} \right)_{fatigue} + \frac{1}{3600f} \left(\frac{da}{dt} \right)_{creep} \quad (26)$$

$$\frac{da}{dt} = \left(\frac{da}{dt} \right)_{fatigue} + \left(\frac{da}{dt} \right)_{creep} = 3600f \left(\frac{da}{dN} \right)_{fatigue} + \left(\frac{da}{dt} \right)_{creep} \quad (27)$$

where da/dt is crack growth rate in mm/hour, da/dN is crack growth rate per cycle (mm/cycle), and f is frequency in Hz. It is well known that the creep damage accumulation is different from damage caused by cyclic loading. According to this a fracture process zone will be different for both creep and fatigue cases.

MATERIAL PROPERTIES

The material used in the tests and numerical calculations is 20CrMoV5 steel, which is used for main power plant components such as steam piping and reheat tubes. All mechanical properties were obtained on smoothed round specimens. Uniaxial tension and fatigue tests on the BISS Nano 25kN servo-hydraulic test system according to ASTM E8, E466 and E646 standards were performed. Creep and damage accumulation law properties were obtained by using UTS-1300-1-50-1-A test system. More details about creep damage accumulation law constants determination



algorithm can be found in [27]. The main mechanical properties of analyzed material were obtained at the elevated temperature of 550⁰ C and summarized in Tab. 1.

Property	Value
Young's modulus, E [GPa]	154
Poisson's ratio, ν	0.3
Strain hardening coefficient, α	2.06
Strain hardening exponent, n_p	8.22
Yield stress, σ_0 [MPa]	265
Ultimate true tensile strength, σ_f [MPa]	577
Norton's constant, B [1/(Pa ⁿ ·hr)]	0.57×10 ⁻²⁴
Norton's constant, n_{cr}	2.52
Damage constant, D [1/(Pa ⁿ ·hr)]	0.54×10 ⁻²⁴
Damage constant, m	2.8
Fatigue strength coefficient, σ'_f [MPa]	750
Fatigue coefficient, c	-0.087

Table 1: The main mechanical properties of 20CrMoV5 steel.

The creep-rupture strength σ_{fr} at a specified temperature T (Kelvin) and creep time t (hours) for 20CrMoV5 steel approximately can be found from the following equation [25]:

$$\sigma_{fr} = (-156.38 \cdot \ln(G) + 586.36) \quad (28)$$

$$G = T(\ln(t) - 2\ln(T) - 24.1) \cdot 10^{-3} \quad (29)$$

EXPERIMENTAL STUDY

One of the main purposes of this present work is to study the influence of the damage accumulation on the creep-fatigue crack growth rate. To this end, subjects for both the experimental studies and numerical analyses are compact tension specimens, which are the most frequently used for characterizing the crack growth rate resistance. In general, the selected dimensions are as recommended by the ASTM test standard [23]. The length of the specimen working area was chosen $w = 50$ mm and specimen thickness $b = 12.5$ mm.

The creep-fatigue tests were carried out on the UTS-110MH-5-0U test system with a high-temperature oven and high-precision crack opening displacement extensometer (Fig. 1a). The potential drop and unloading compliance methods were used to monitor the crack length during the tests. The waveforms for the loading and unloading portions were trapezoidal, and the loading/unloading times were maintained constant (5 s rise and decay times). A hold time of a predetermined duration, namely 60 s, was superimposed on the trapezoidal waveforms at the maximum load. Parts of the 5 s loading and unloading cycles represent low-cycle deformation, while the 60 s dwell under a constant load represents creep (Fig. 1b).

The tests were conducted at 550°C on special high-temperature test equipment with a load ratio R of 0.1 and maximum load value P_{max} of 11 kN. The specimens were pre-cracked to an initial crack length-to-width ratio a/w of approximately 0.36 to 0.39 under cyclic loading at room temperature. For each tested specimen during the creep-fatigue tests the



amplitude ratio was changed from $R=0.1$ to $R=0.9$ two times for a short time period. This practice helped to determine the intermediate crack front positions during the creep–fatigue loading between the initial (pre-crack) and the final crack fronts (Fig.2).

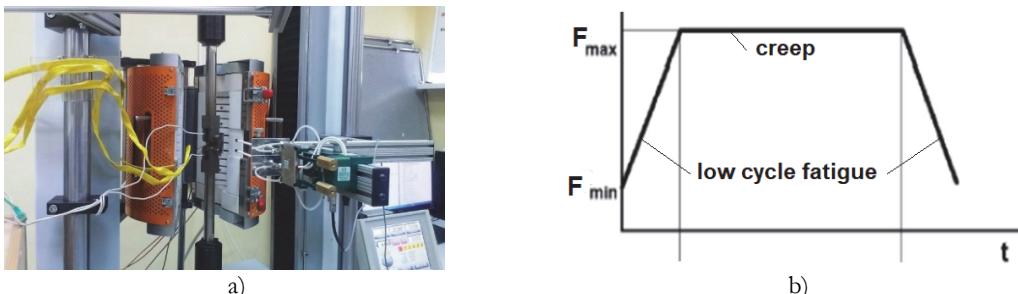


Figure 1: Tests equipment (a) and the load waveform (b).

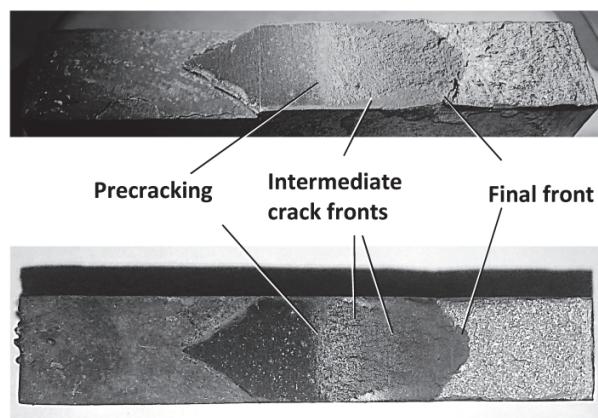


Figure 2: Fracture surfaces.

After total failure of the specimen, measurements of the crack sizes were taken for four positions of the crack front by means of an optical microscope. For each front, the crack size at five equally spaced points centered on the specimen mid-thickness line was measured along the crack front.

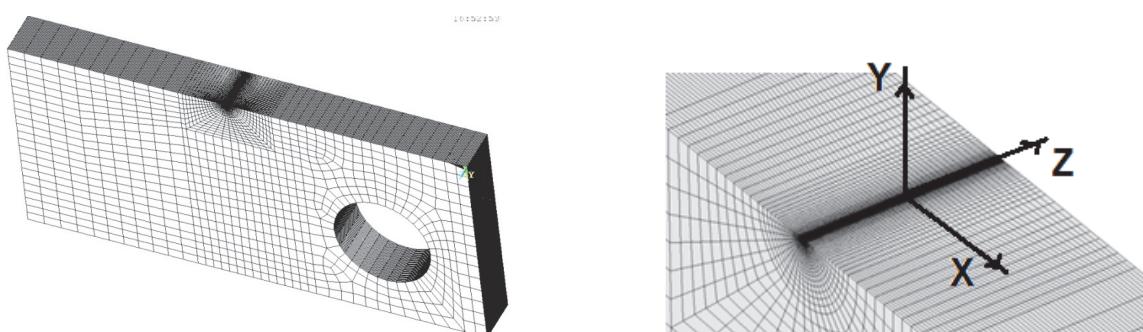


Figure 3: Finite element method meshes for compact tension specimen.

FINITE ELEMENT MODEL

Full-field finite element analysis is performed using ANSYS software to study the crack-front stress fields for compact tension specimen. Along the thickness direction, an identical planar mesh is repeated from the symmetry plane to the free surface. To catch the drastic change of the stress field near the free surface, the thickness of

successive finite element layers is exponentially reduced from the mid-plane toward the free surface. The radial sizes of the finite elements are varied according to the geometric progression. According to the symmetrical properties, one quarter of the actual structure was selected to establish the three-dimensional finite element model. The twenty-node quadrilateral brick isoparametric three-dimensional solid elements were used to model for the specimen configuration. Typical finite element meshes for the compact tension specimen are illustrated in Fig. 3. Programming of the numerical calculations includes the analysis for the specified experimental combinations of creep time/number of fatigue cycles and crack length in a compact specimen. Three position of the crack front were considered ($a/w = 0.38, 0.55$ and 0.7). For $a/w = 0.55$ and 0.7 a curvilinear crack fronts were simulated.

RESULTS AND DISCUSSION

As result of a finite element simulations a stress-strain fields were obtained for different combinations of crack length and creep time. The governing parameter behavior of crack tip fields I for the compact specimen is plotted in Fig. 4 along the crack front towards the thickness direction.

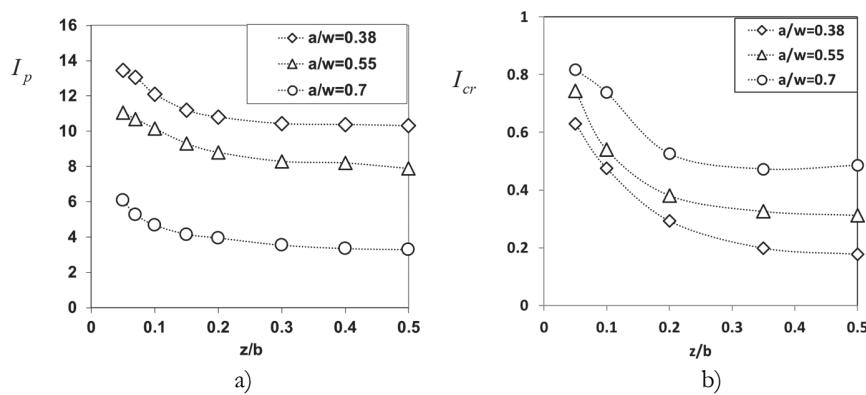


Figure 4: I -factor distributions along crack front for plastic (a) and creep (b) state.

It should be noted that in real solids, the governing parameter of the crack tip field the I -integral is very sensitive to variations of the material properties, load level and geometry. These distributions of the In-integral are used to calculate the nonlinear stress intensity factors in the compact tension specimen.

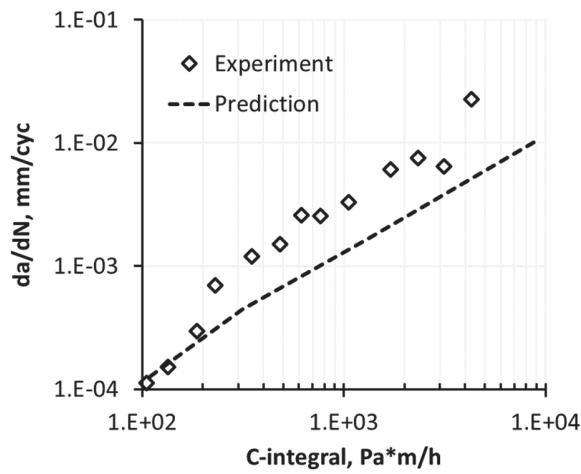


Figure 5: Comparison of crack growth rate prediction with experimental results for steel 20CrMoV5.

The main objective of the creep-fatigue tests is to provide substantiation for the theoretical models. Fig. 5 illustrates the prediction of the creep-fatigue crack growth rate for the 20Cr1MoV5 steel plotted against C^* -integral for the compact tension specimen. The experimental C^* -integral values were obtained from the crack length, determined using the



technique according to ASTM E2760 [23]. The experimental data of all tested specimens fall within a relatively narrow scatter band. A good agreement with the theoretical predictions is observed. By the author opinion the shift between prediction and experiment is caused by default value of initial damage w_0 for each finite element model. In real specimens the value of initial damage during creep-fatigue depends on load history. For short cracks and small creep times it can be neglected, but influence over time enhanced. The proposed models of the crack growth rate are formulated in terms of nonlinear stress intensity factors. This model can describe the behavior a wide range of the material different properties for creep-fatigue interaction.

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